



Probing R-Symmetric SUSY Lepton Flavor

@Project X

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- *The SUSY lepton flavor problem*
- *The Minimal R-Symmetric Supersymmetric Model (MRSSM)*
Phys.Rev.D78:055010,2008
[Graham D. Kribs, Erich Poppitz, Neal Weiner](#)
- μ to e flavor constrains on the MRSSM
Phys.Rev.D82:035010,2010
[Ricky Fok, Graham D. Kribs](#)
- *Improving the current constrains at Project X*



The SUSY Lepton Flavor Problem



SUSY lepton flavor problem



In the MSSM

m^2	m_{ij}^2	m_{ij}^2	m_{ij}^2	m_{ij}^2	m_{ij}^2
m_{ij}^2	m^2	m_{ij}^2	m_{ij}^2	m_{ij}^2	m_{ij}^2
m_{ij}^2	m_{ij}^2	m^2	m_{ij}^2	m_{ij}^2	m_{ij}^2
m_{ij}^2	m_{ij}^2	m_{ij}^2	m^2	m_{ij}^2	m_{ij}^2
m_{ij}^2	m_{ij}^2	m_{ij}^2	m_{ij}^2	m^2	m_{ij}^2
m_{ij}^2	m_{ij}^2	m_{ij}^2	m_{ij}^2	m_{ij}^2	m^2

L

R

The m_{ij}^2 are tightly constrained!

Nucl.Phys.B783:112–142,2007

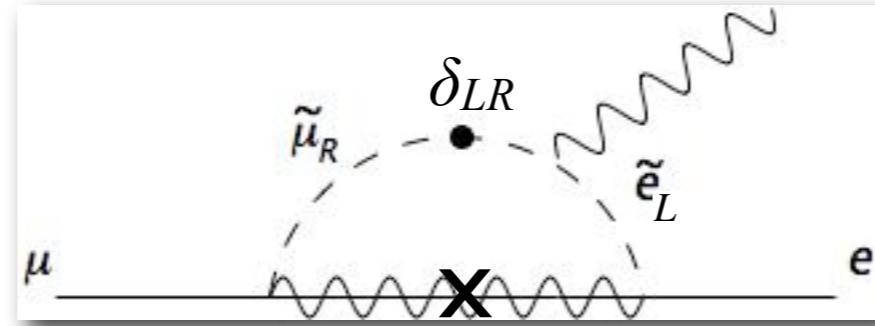
M. Ciuchini, A. Masiero, P. Paradisi,
L. Silvestrini, S.K. Vempati, O. Vives

$$\delta_{ij} \equiv \frac{m_{ij}^2}{m^2} < 10^{-3} \text{ in } mSUGRA$$

WHY?? This is the SUSY lepton flavor problem

In the MSSM, the most constraining process is

$$\mu \rightarrow e\gamma$$



$$\delta_{LR} \equiv \frac{m_{LR}^2}{m^2}$$

$$\delta_{LR} < 3 \times 10^{-5}, \text{ mSUGRA}$$

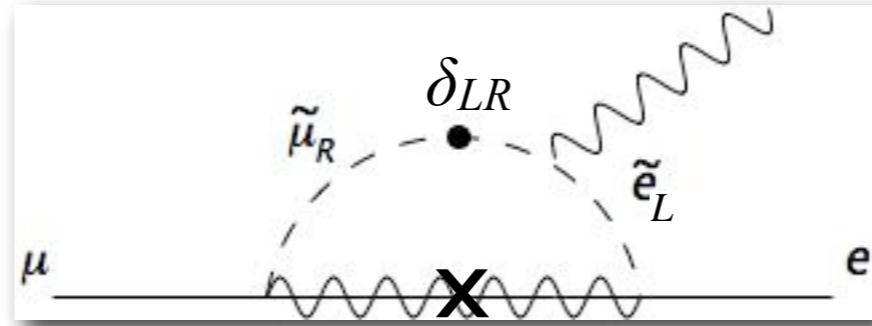
It contains a **left-right mixing** of the sleptons
and a **chirality flip** of the Gaugino/Higgsino

SUSY lepton flavor problem



In the MSSM, the most constraining process is

$$\mu \rightarrow e\gamma$$



$$\delta_{LR} < 3 \times 10^{-5}, \text{ mSUGRA}$$

As we shall see, this is diagram **NOT allowed** in the MRSSM

Current limits on m_{ij}^2 is much more relaxed, potentially
solving the lepton flavor problem

Project X will be able to determine whether the MRSSM is a
solution



The Minimal R-Symmetric Supersymmetric Model **(MRSSM)**

Phys.Rev.D78:055010,2008

Graham D. Kribs, Erich Poppitz, Neal Weiner



Extending R-parity



Start with the MSSM with a Z_2 R-parity

SM fields	R Parity
Leptons (L,R)	Even
Gauge Bosons	Even
Higgs scalar	Even

SUSY fields	R Parity
Sleptons (L,R)	(Odd, Odd)
Gauginos	Odd
Higgsinos	Odd

Now Extend the R-parity to a $U(1)$ continuous symmetry
i.e. conserved R-charges for each field



Extending R-parity



The MSSM with a $U(1)$ R-Symmetry

SM fields	R Charge
Leptons (L,R)	0
Gauge Bosons	0
Higgs scalar	0

SUSY fields	R Charge
Sleptons (L,R)	(-1, 1)
Gauginos	1
Higgsinos	-1



Extending R-parity



The MSSM with a $U(1)$ R-Symmetry

SM fields	R Charge
Leptons (L,R)	0
Gauge Bosons	0
Higgs scalar	0

SUSY fields	R Charge
Sleptons (L,R)	(-1, 1)
Gauginos	1
Higgsinos	-1

Consequences

- a) *Gaugino / Higgsino majorana masses terms violates R charge by 2.*

Solution: Make gauginos and higgsinos Dirac.

- b) *Left-right slepton mixing also violates R charge by 2.*

Gaugino and Higgsino masses

Introduce R-partners of Gauginos and Higgsinos

MRSSM fields	R Charge
R-Gauginos	-1
R-Higgsinos	1

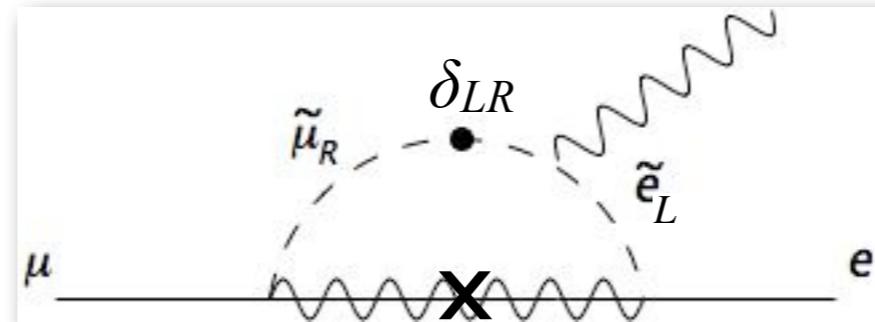
SUSY fields	R Charge
Sleptons (L,R)	(-1, 1)
Gauginos	1
Higgsinos	-1

This allows for Dirac masses of inos.

The R-gaugino and gauginos are the right and left chiralities of the Dirac Gauginos

These new fields are not important in LFV processes and will not be considered further.

Recall the most constraining diagram in the MSSM is



This diagram vanishes in the MRSSM, two reasons:

- a) no L-R mixing, $\delta_{LR} = 0$.
- b) The chirality flip on the ino line flips the ino to its R-partner, which does not couple to leptons.

Question: What are the constraints on the slepton mixing in the MRSSM?? What is the dominant process?



μ to e flavor mixing constraints on the MRSSM

Phys.Rev.D82:035010,2010
[Ricky Fok, Graham D. Kribs](#)



The Sensitive Parameters



Bino mass M_1

H_d mass μ_d

Slepton masses m_1 m_2 ($m_2 = 1.5 m_1$)

$\tilde{\mu}$ \tilde{e} mixing $\sin 2\theta_L$ $\sin 2\theta_R$

m^2	m_{ij}^2	0	0	0	0
m_{ij}^2	m^2	0	0	0	0
0	0	m^2	0	0	0
0	0	0	m^2	m_{ij}^2	0
0	0	0	m_{ij}^2	m^2	0
0	0	0	0	0	m^2

L

R

*LR mixing
forbidden in
MRSSM*

L

R

We computed branching ratios of the following processes

$$\mu \rightarrow e\gamma$$

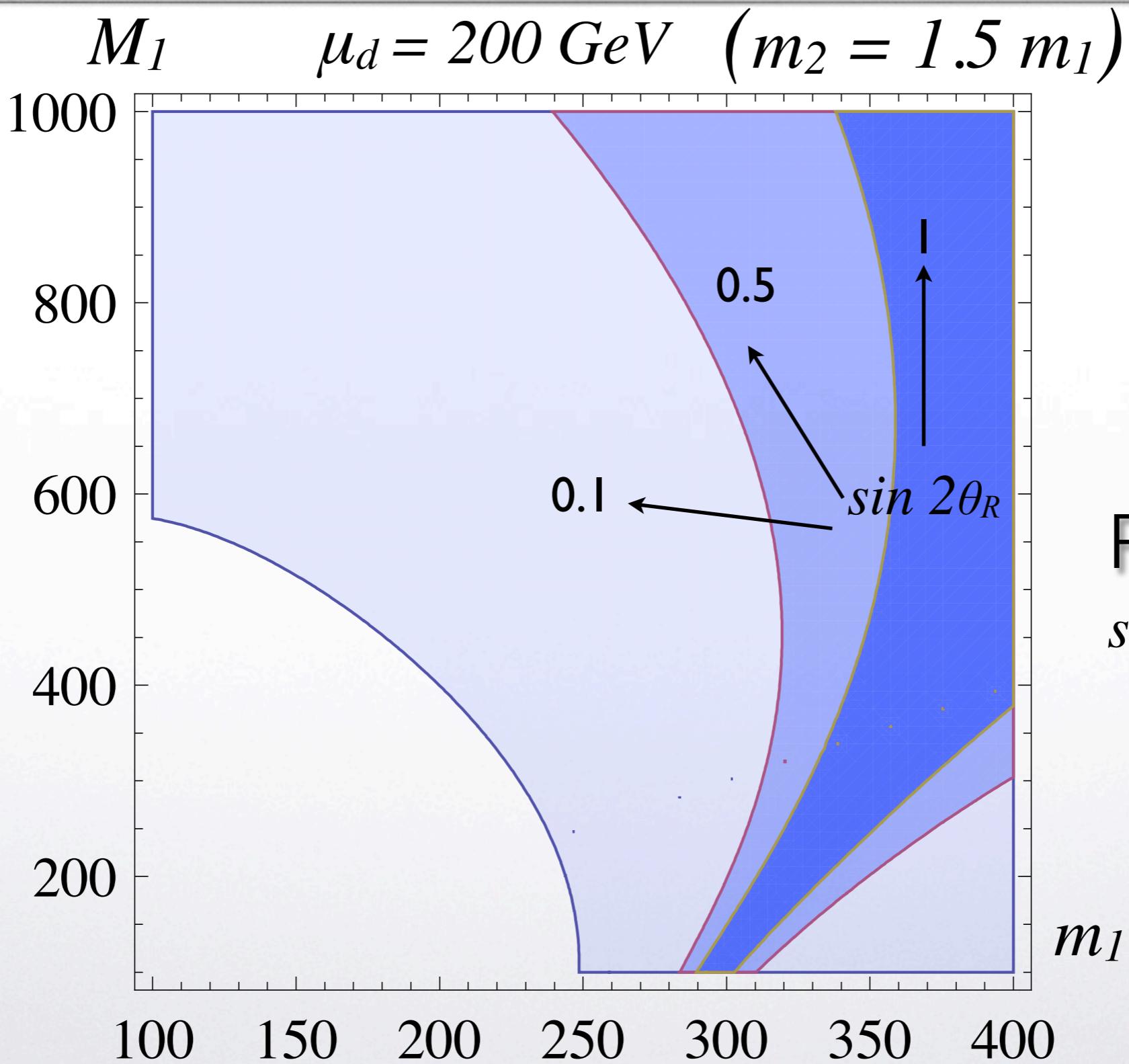
$\mu \rightarrow e$ conversion
in gold

$$\mu \rightarrow eee$$

With only either one of θ_L or θ_R to be non-zero

Set $\tan \beta = 3$, and set the heavier slepton mass to be $m_2 = 1.5m_1$

Then we set limits on the parameter space
 $(m_B, \mu_d, m_1, \sin 2\theta_{L,R})$ with current experimental bounds



RH
 $\sin 2\theta_L = 0$

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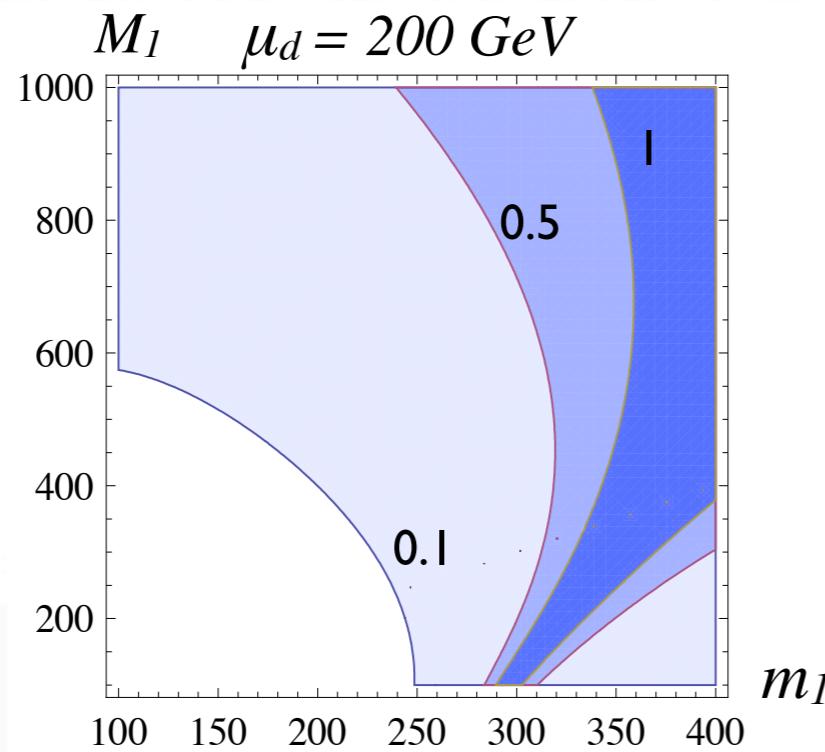


Right-handed mixing

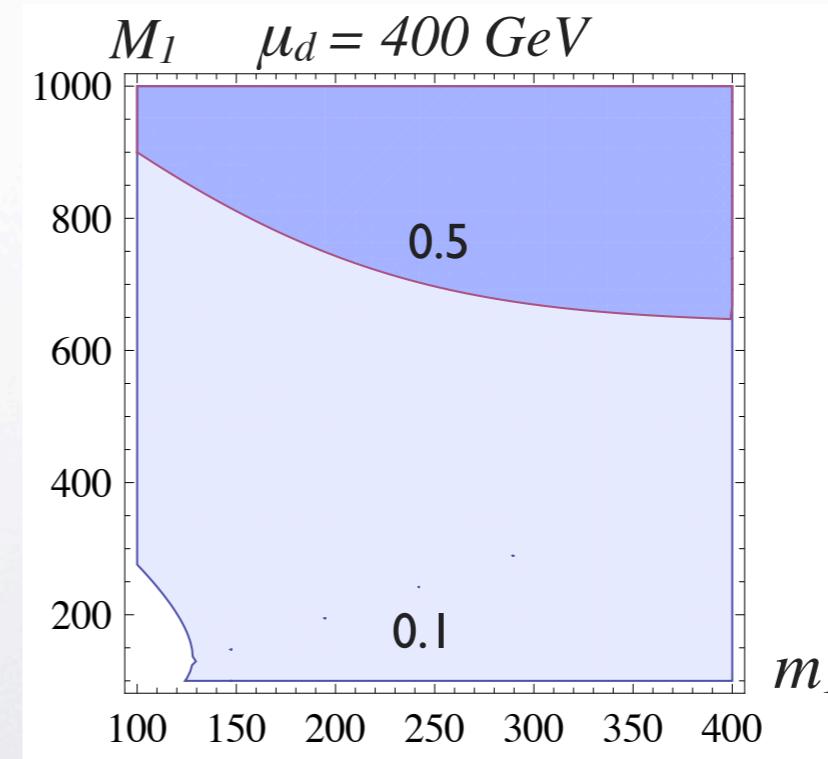
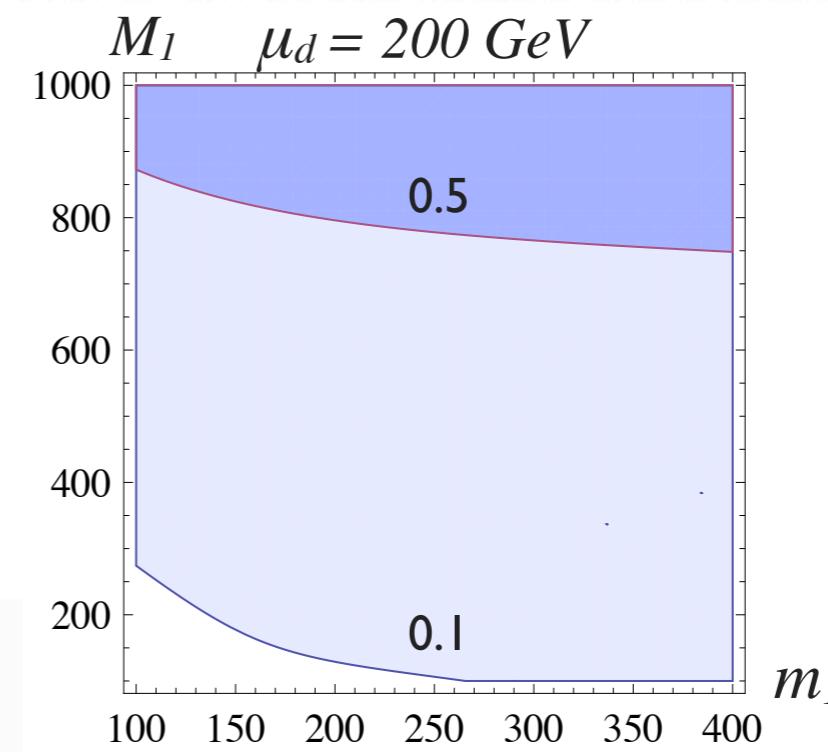
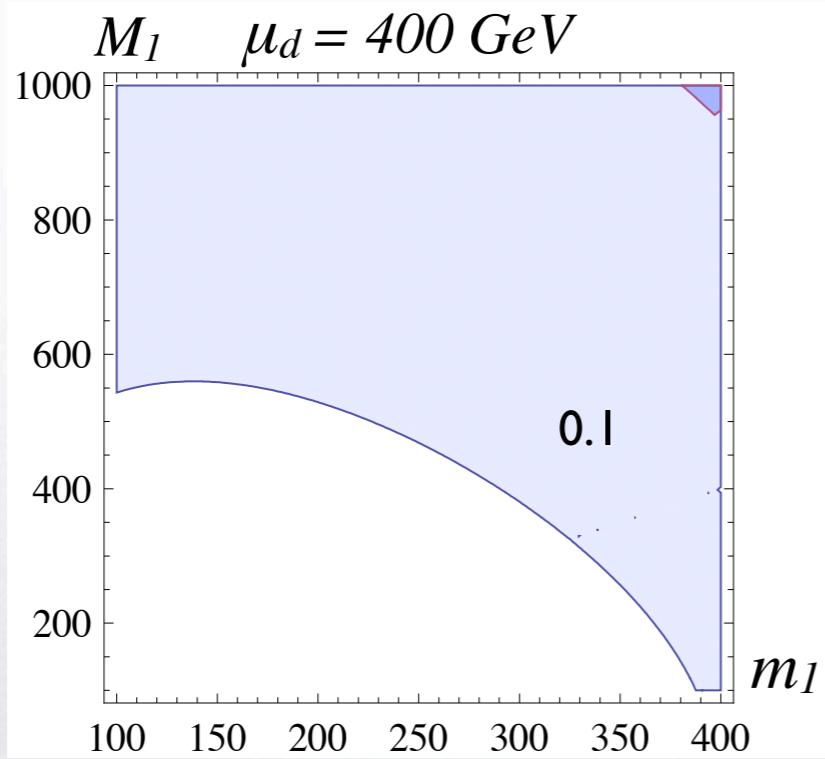
$$\sin 2\theta_L = 0$$



$\mu \rightarrow e\gamma$



$\mu \rightarrow e\gamma$



$(\mu \rightarrow e)_{\text{Au}}$

$(\mu \rightarrow e)_{\text{Au}}$

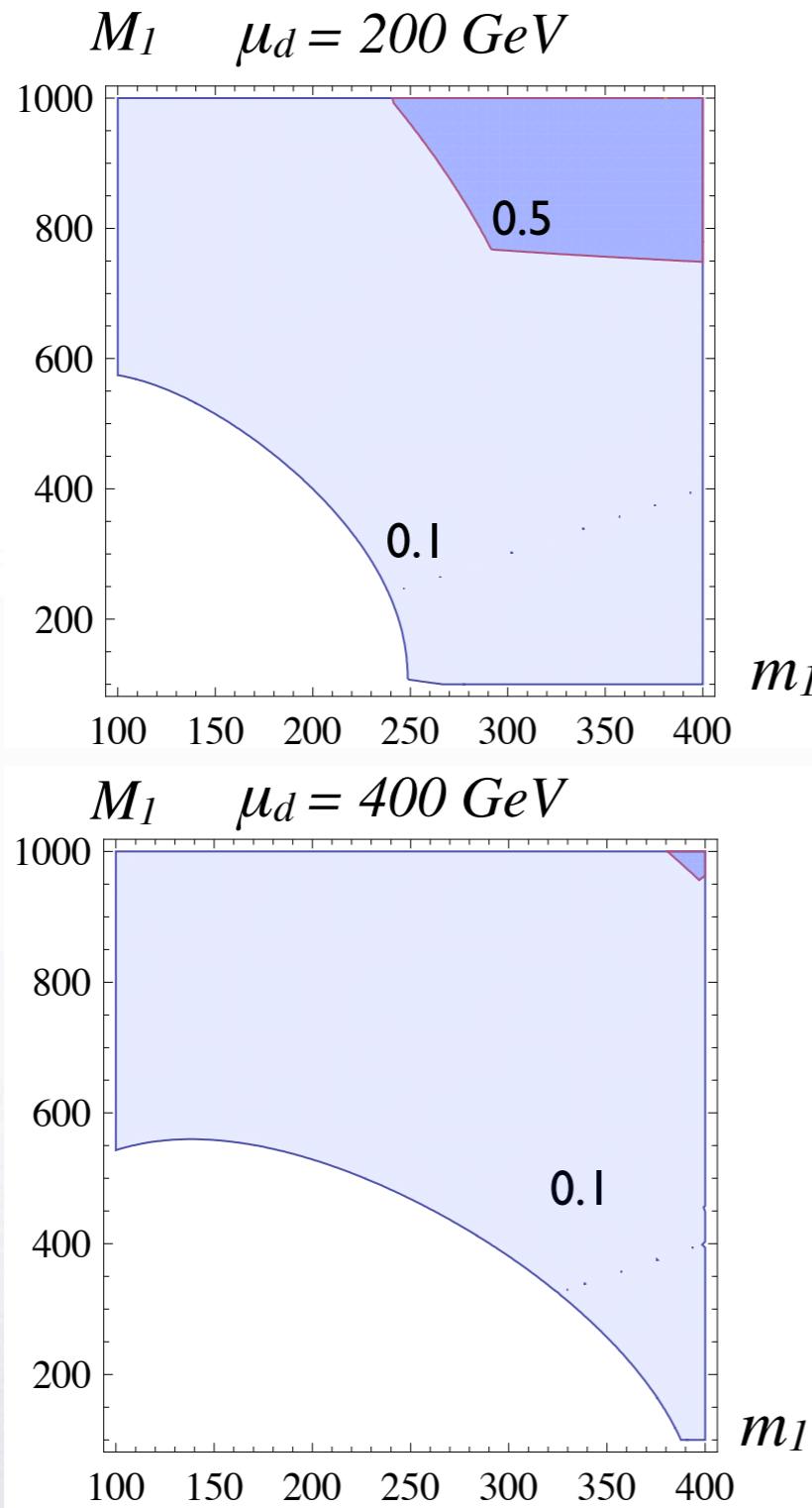
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$\text{BR} < 2.4 \times 10^{-12}, \text{MEG}$

$\text{BR} < 7.0 \times 10^{-13} \text{ SINDRUM II}$



Combined Right handed mixing



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$O(0.1)$ mixing allowed

$\mu \rightarrow eee$ gives looser bounds

$\mu \rightarrow e\gamma$ gives
slightly better bounds

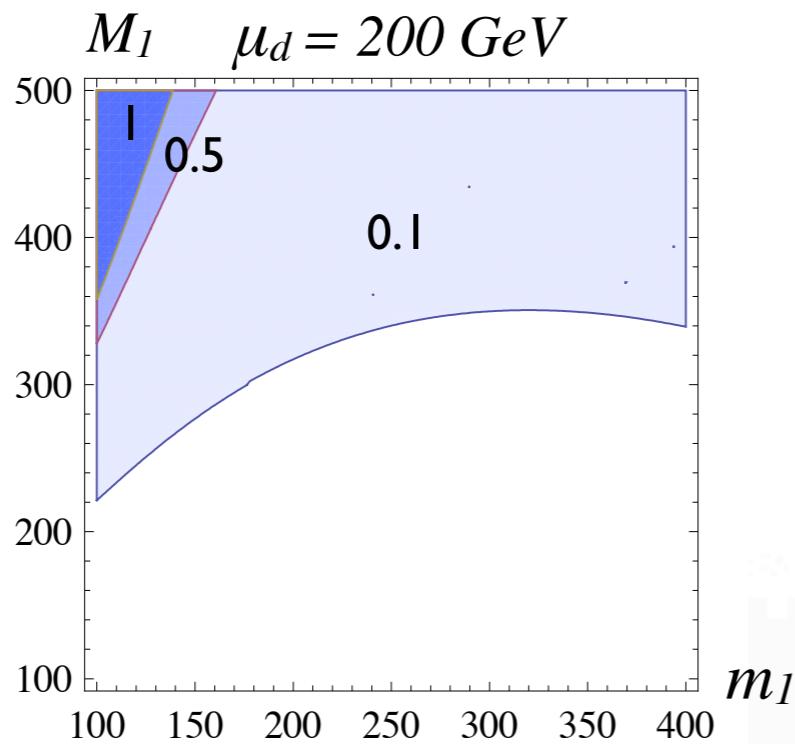


Left-handed mixing

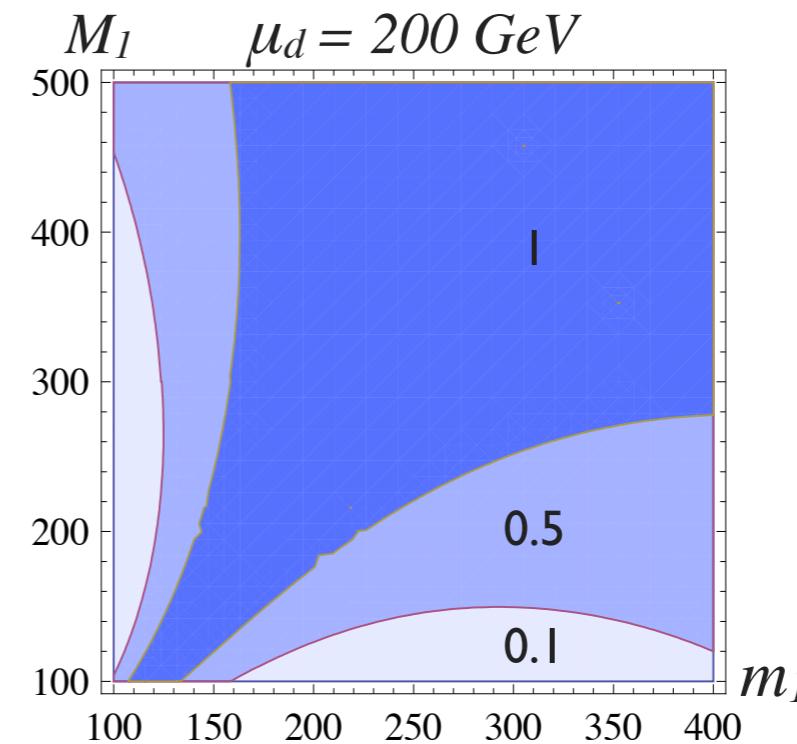
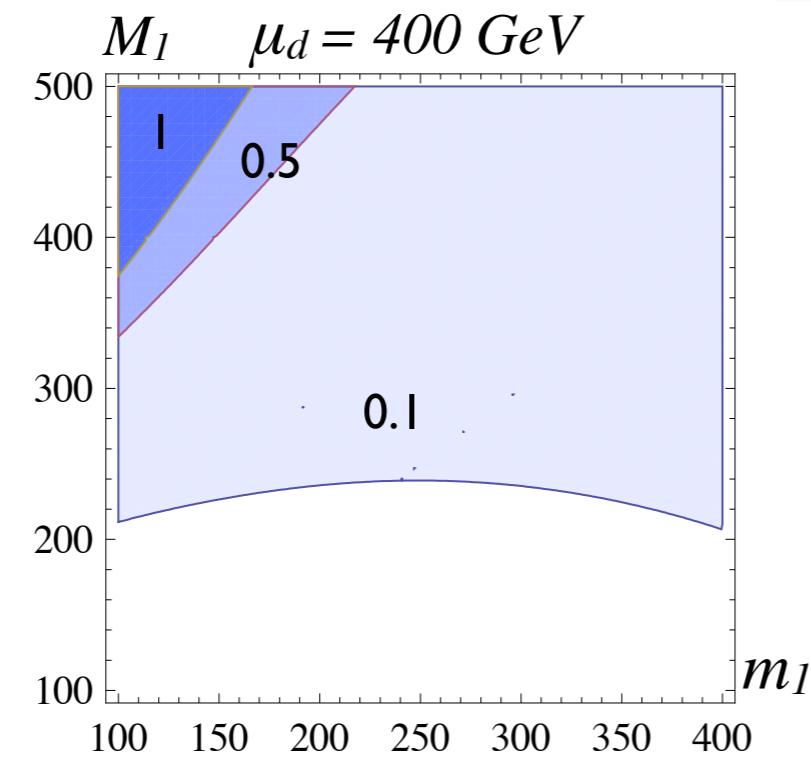
$$\sin 2\theta_R = 0$$



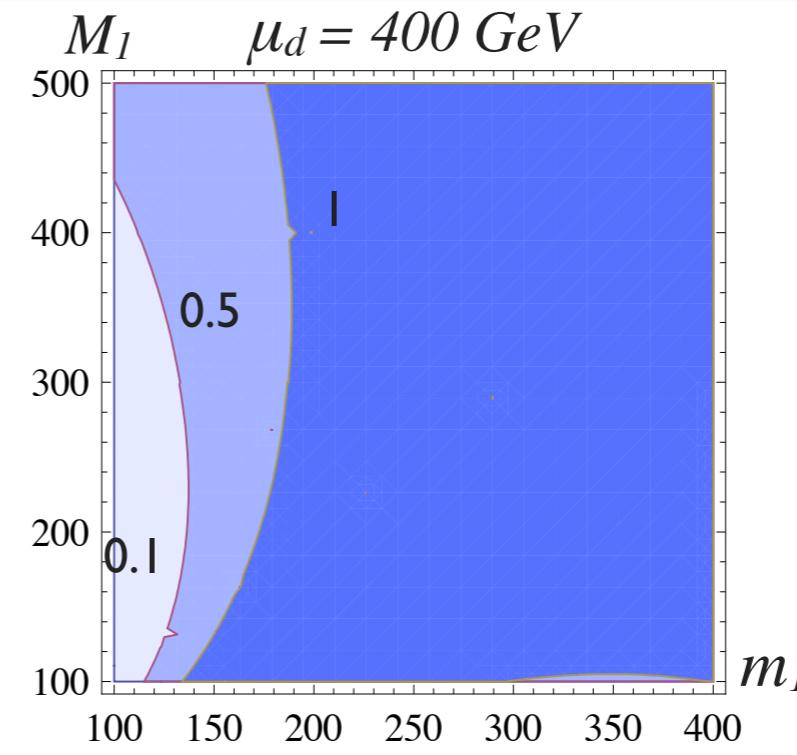
$\mu \rightarrow e\gamma$



$\mu \rightarrow e\gamma$



$(\mu \rightarrow e)_\text{Au}$



$(\mu \rightarrow e)_\text{Au}$

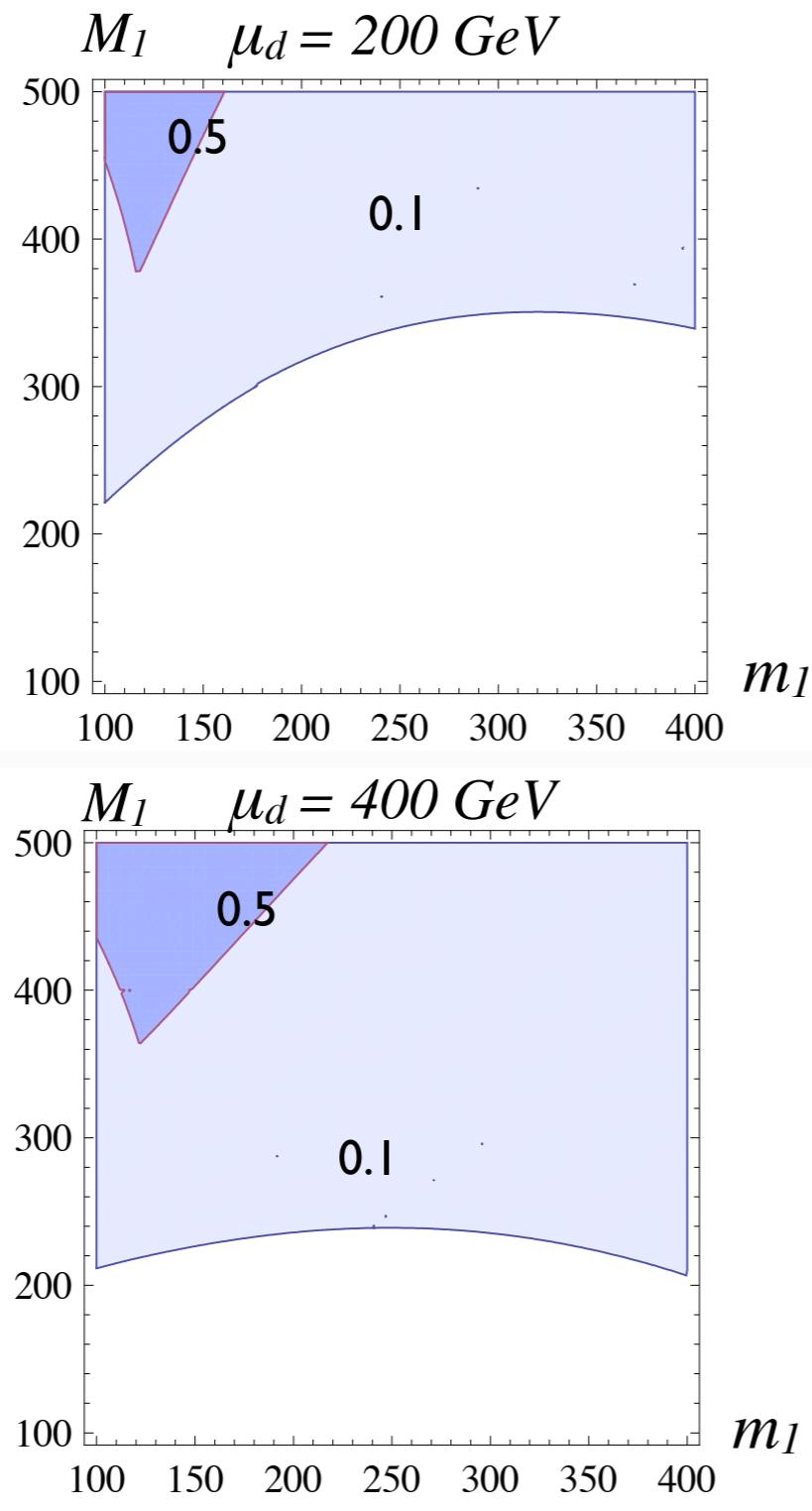
Fok,
Kribs

$\text{BR} < 2.4 \times 10^{-12}$, MEG

$\text{BR} < 7.0 \times 10^{-13}$ SINDRUM II



Combined Left handed mixing



$O(0.1)$ mixing allowed

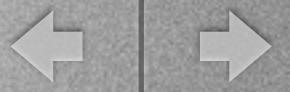
$\mu \rightarrow eee$ gives looser bounds

$\mu \rightarrow e\gamma$ gives
slightly better bounds

Fok,
Kribs



What does the plots mean?



$\mu \rightarrow e$ conversion from right handed mixing is coherent,
can be used to constrain $\sin 2\theta_R$

Current bounds allow $\sin 2\theta_{L,R}$ of $O(0.1)$ for
most masses in the sub-TeV range

This does **NOT** say anything definite about the
MRSSM as a solution for the lepton flavor
problem

But Project X can!



Probing MRSSM LFV with Project X



Expected limits



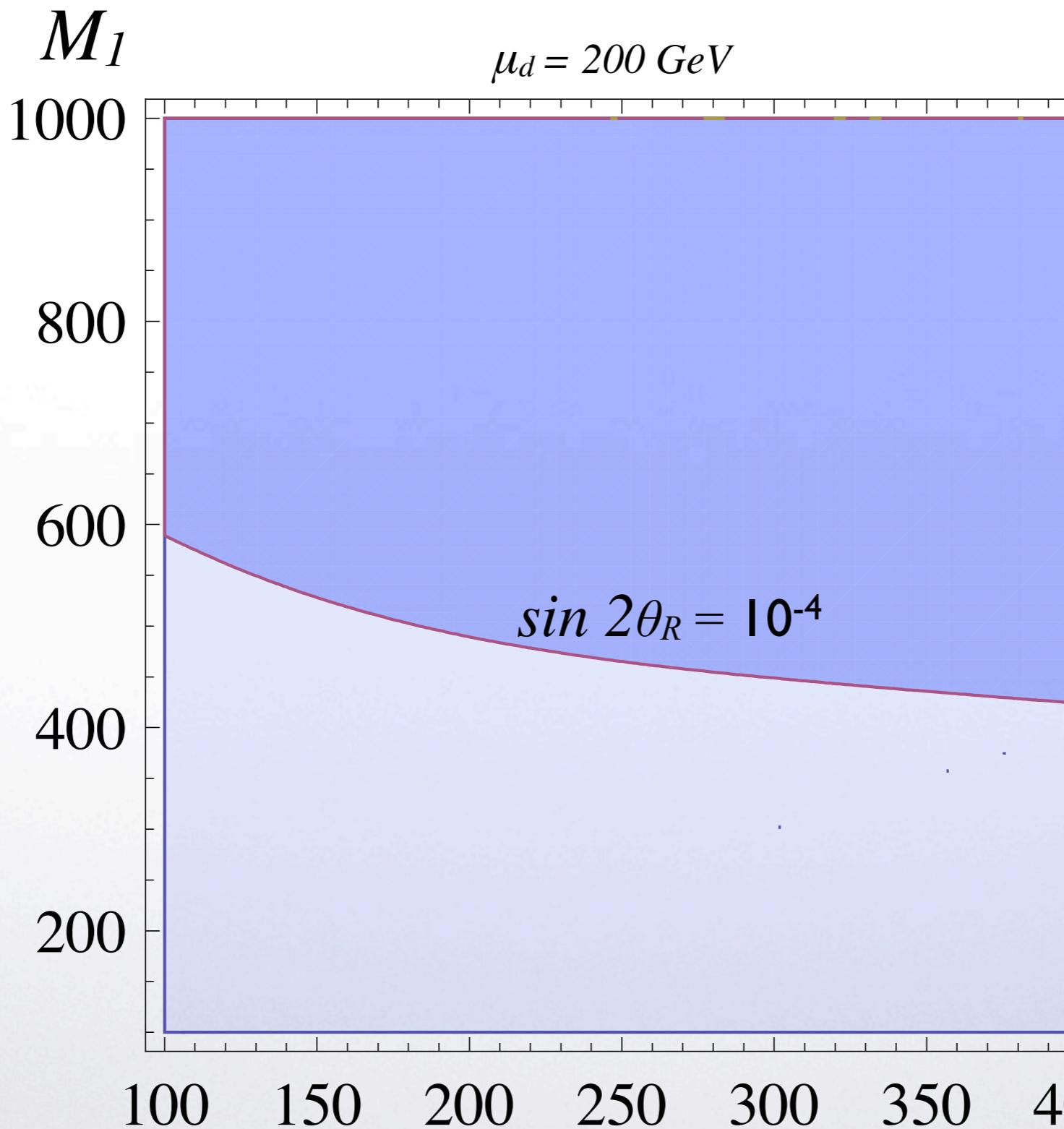
Expected sensitivities at Project X

	Improvement in sensitivity	Expected limits on $\sin 2\theta$
$\mu \rightarrow e$ conversion	10^6	$10^{-3} - 10^{-4}$
$\mu \rightarrow eee$	$10^3 - 10^4$	10^{-3}
$\mu \rightarrow e\gamma$???	???

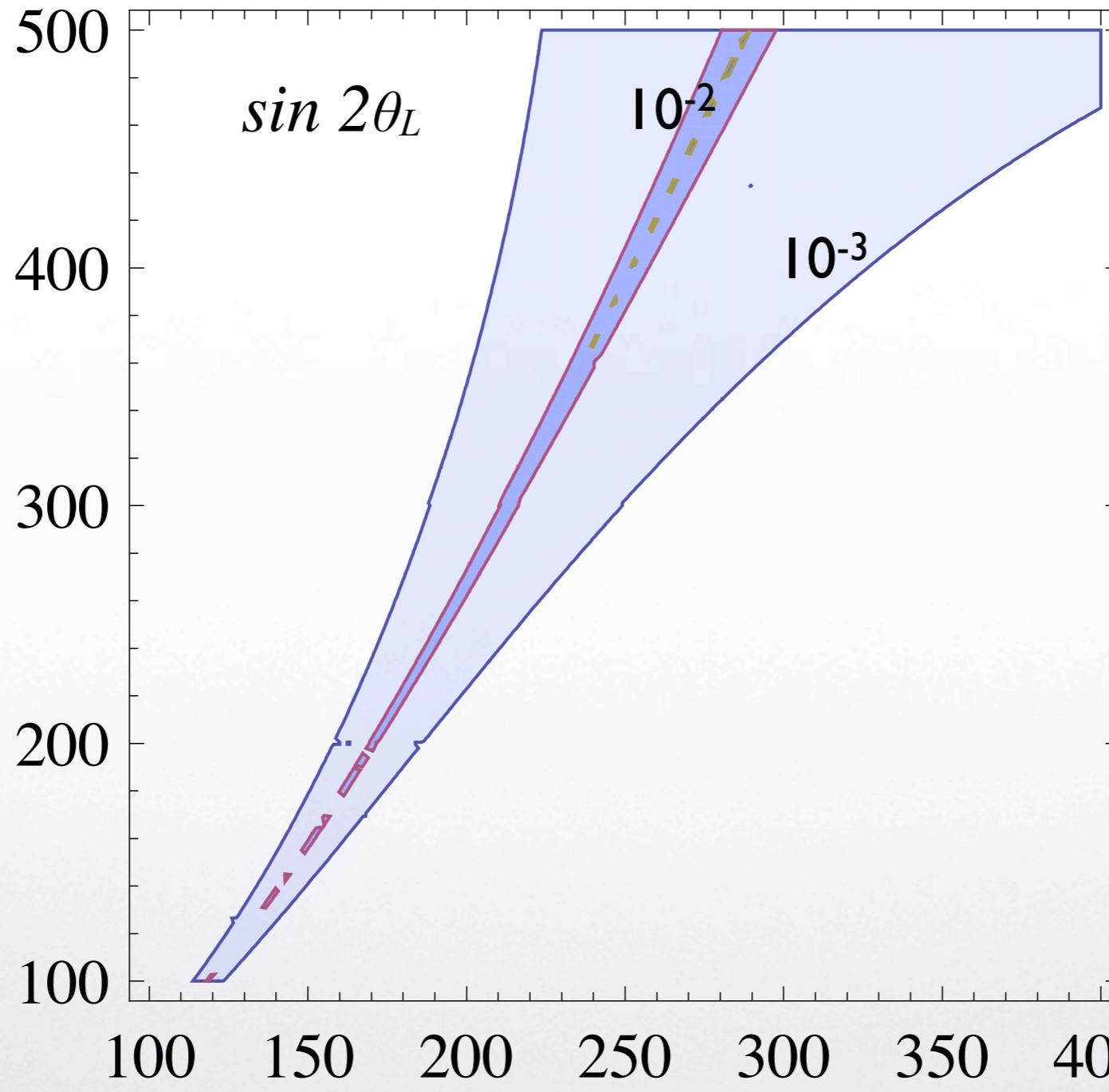
Recall that $\mu \rightarrow e$ conversion due to RH slepton mixing is not affected by destructive interference

Checked numerically that $\mu \rightarrow eee$ does not change on the expected limits from $\mu \rightarrow e$ conversion significantly

Project X and the Science of the Intensity Frontier, white paper, Fermilab, 2010.
<http://www.fnal.gov/pub/projectx/pdfs/ProjectXwhitepaperJan.v2.pdf>



Fok

M_1 $\mu_d = 200 \text{ GeV}$  $LH \text{ mixing}$ m_1

Fok



Expected limits



m^2	m_{ij}^2	0	0	0	0
m_{ij}^2	m^2	0	0	0	0
0	0	m^2	0	0	0
0	0	0	m^2	m_{ij}^2	0
0	0	0	m_{ij}^2	m^2	0
0	0	0	0	0	m^2

L R

Project X can constrain these elements down to 10^{-4} , assuming $O(1)$ slepton mass splitting



Expected limits



m^2	m_{ij}^2	0	0	0	0
m_{ij}^2	m^2	0	0	0	0
0	0	m^2	0	0	0
0	0	0	m^2	m_{ij}^2	0
0	0	0	m_{ij}^2	m^2	0
0	0	0	0	0	m^2

L R

Project X can constrain these elements down to 10^{-3} , assuming $O(1)$ slepton mass splitting

If Project X does not find any $\mu \rightarrow e$ conversion events, the flavor problem persists and cannot be explained by the MRSSM

Conclusion



*Search for rare μ decays @ Project X can **improve limits** on LFV parameters by **orders or magnitude***

*This is demonstrated in the μ_R - e_R mixing in the MRSSM, where the limit on their mixing angle can be **improved by at least 10^3** from the current bound at SUNDRUM II.*

*Unless the two sleptons are degenerate, or they mostly mix with the stau. Project X can give a **definite statement** on whether the lepton flavor problem is solved by the MRSSM.*

BACK UP



Insensitive Parameters



All parameters related to the stau:

- *stau masses, mixing with selectron and smuon*

Wino mass

- *Must be of $O(\text{TeV})$ to avoid excessive contribution to electroweak precision parameters.*

H_u masses

- *Charged leptons do not couple directly to H_u*

$\tan \beta$

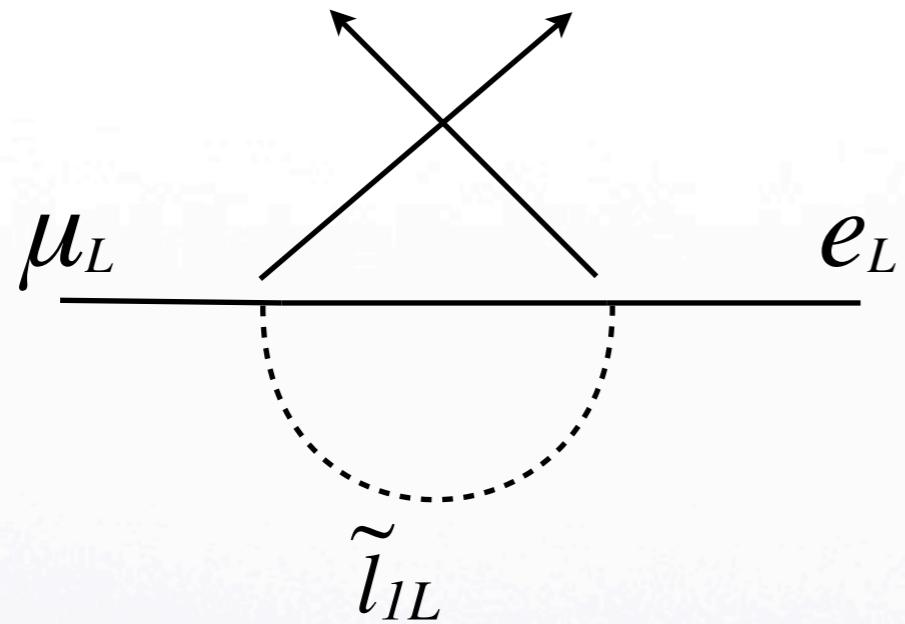
- *Cancels between Yukawa and H_d - B mixing*



Mixing angles



$$\tilde{l}_{IL} = \cos\theta_L \tilde{e}_L + \sin\theta_L \tilde{\mu}_L$$



mixing parameter
 $0 < \sin 2\theta_L < 1$

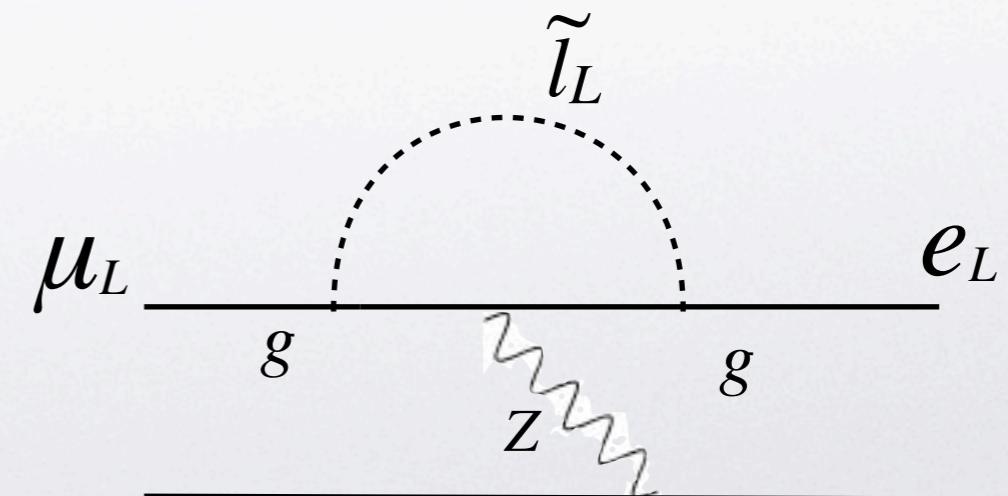
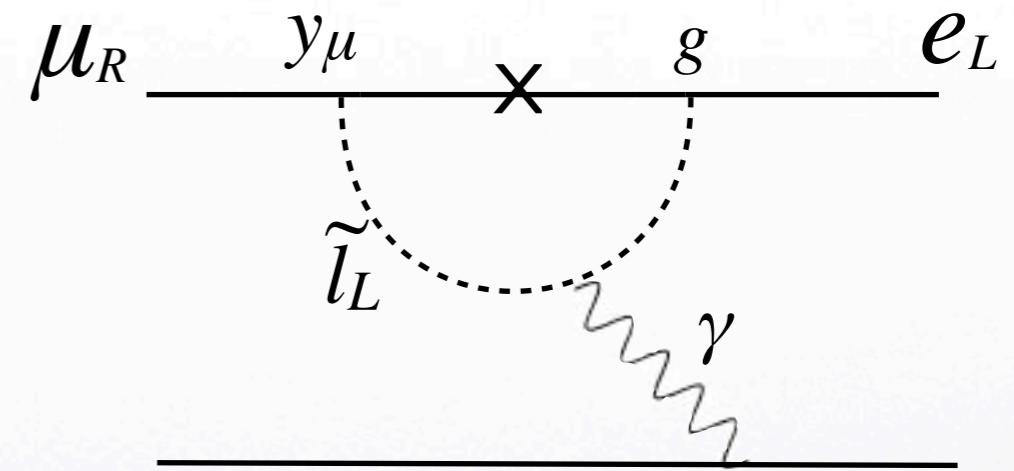
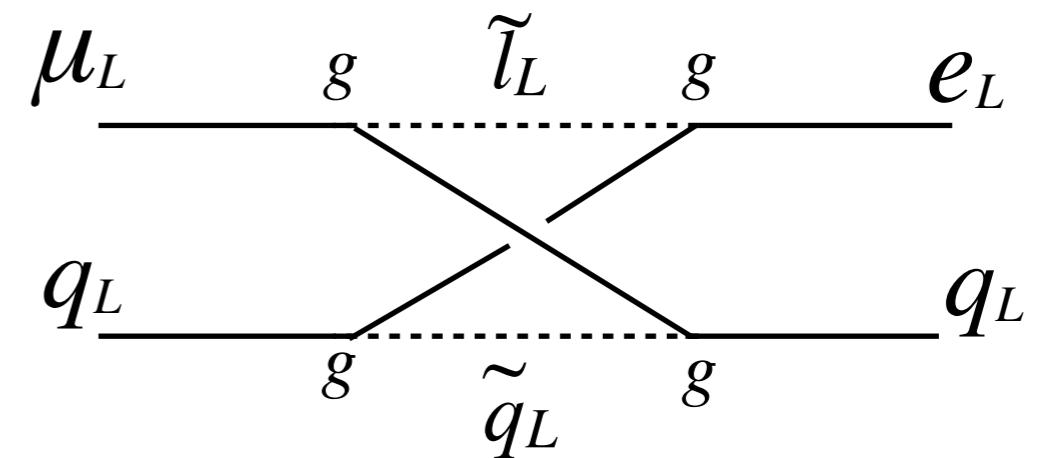
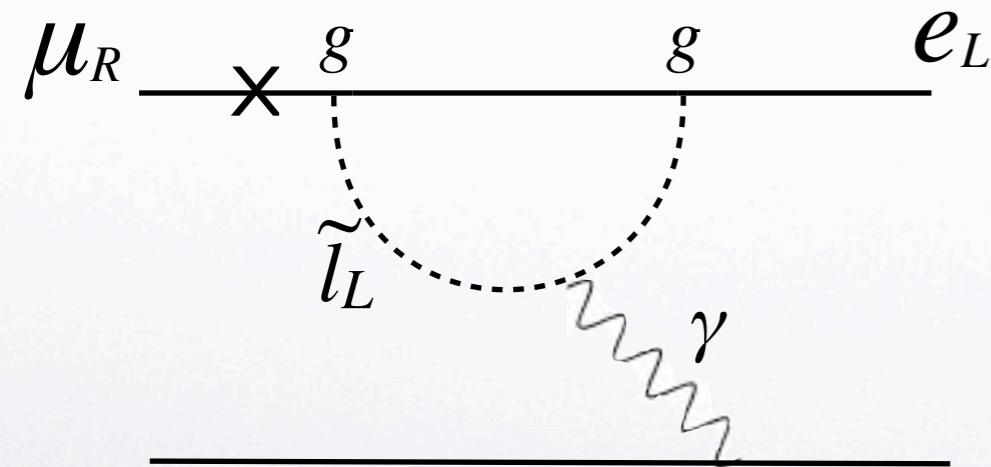
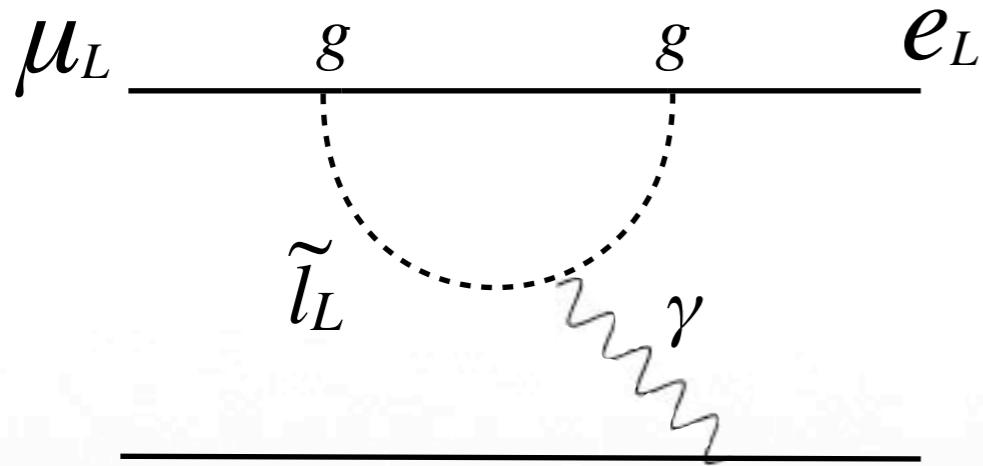
	\tilde{e}_L	$\tilde{\mu}_L$	$\tilde{\tau}_L$			
\tilde{l}_{1L}	c_L	s_L	0	0	0	0
\tilde{l}_{2L}	$-s_L$	c_L	0	0	0	0
	0	0	1	0	0	0
	0	0	0	c_R	s_R	0
	0	0	0	$-s_R$	c_R	0
	0	0	0	0	0	I
	\tilde{e}_R	$\tilde{\mu}_R$	$\tilde{\tau}_R$			



$\mu \rightarrow e$ conversion



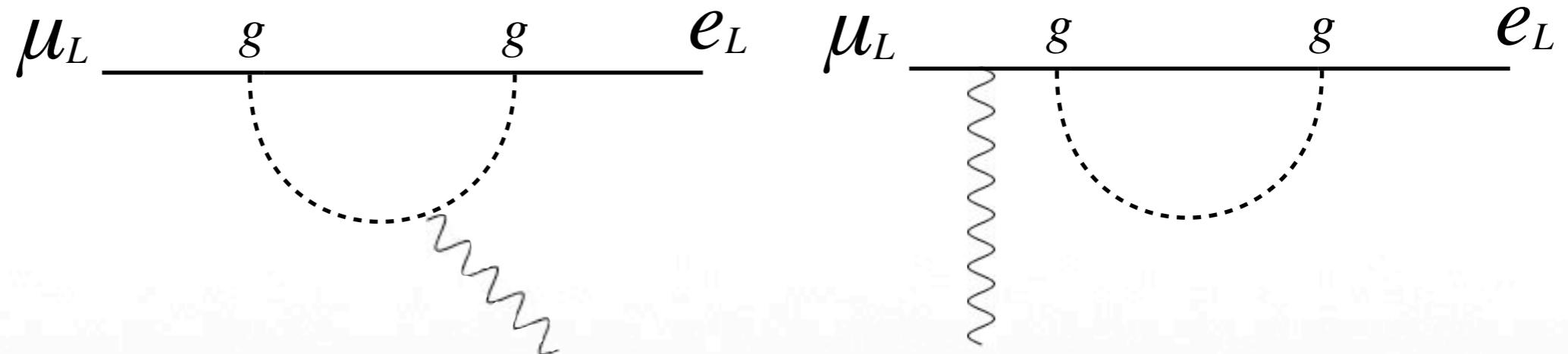
All the relevant diagrams are



Why are chirality flips relevant?!



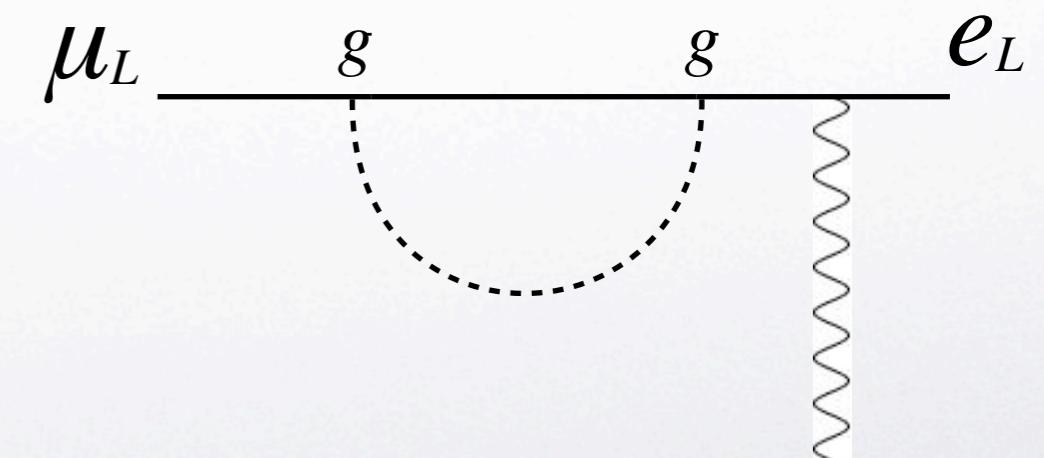
γ Penguin



$$\delta F_I(k^2) - \delta F_I(0), \text{ where } k^2 \sim -m_\mu^2 \\ = \delta F_I'(0)k^2$$

$$\sim \mathcal{O}(\delta F_I) \frac{m_\mu^2}{m_l^2}, \text{ dim 6!}$$

\sim **Same order** as the chirality flipping diagram $m\bar{\psi}\sigma^{\mu\nu}\psi F_{\mu\nu}$

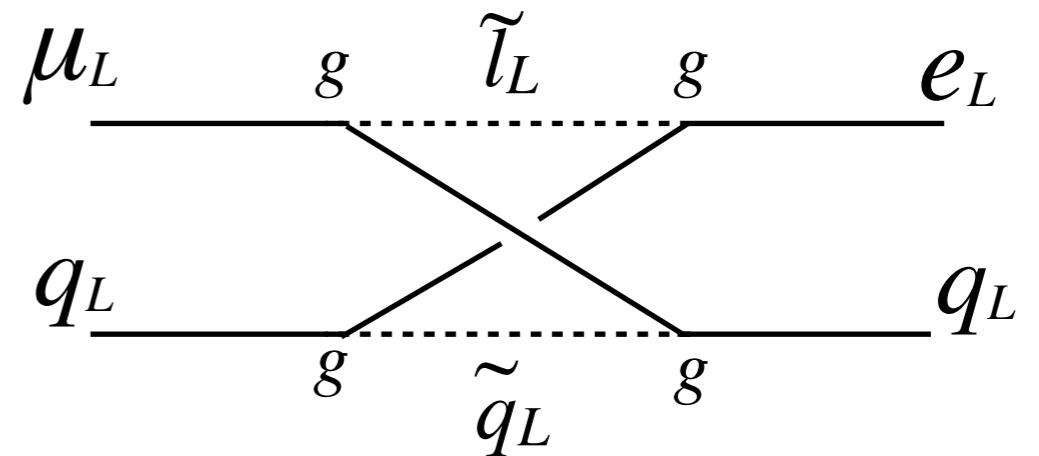




Box

No chirality flip

$$\bar{\psi} \gamma^\mu P_L \psi$$



R-symmetry:

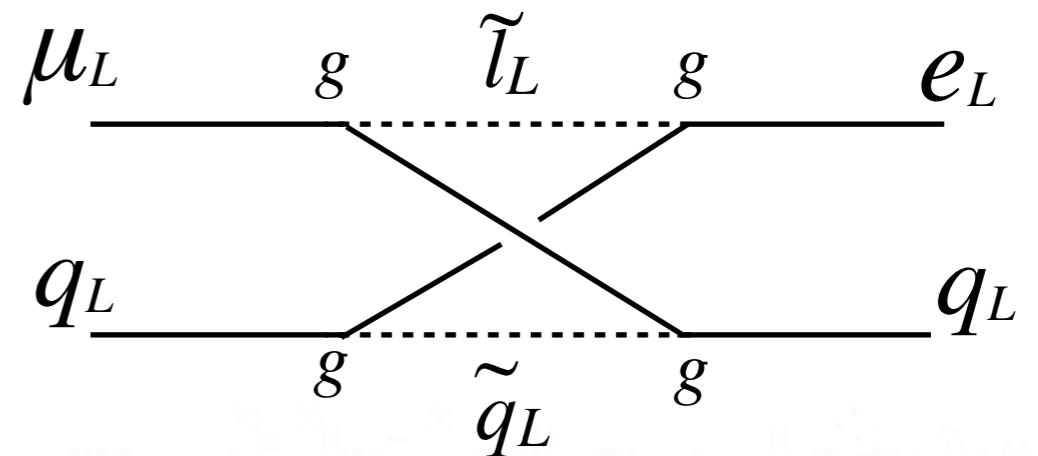
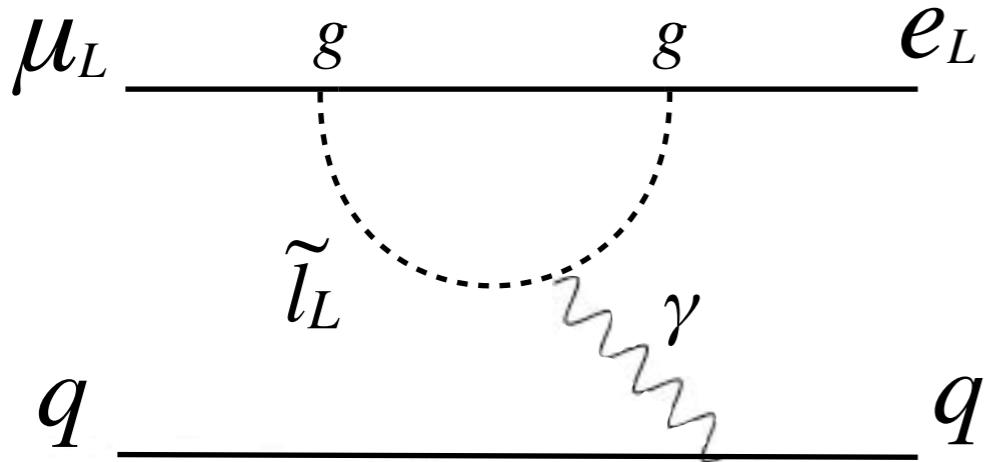
No L-R mixing between sleptons

The **quark line must have the same chirality as the muon!!**

$$(\bar{\psi} \gamma^\mu P_L \psi) (\bar{q} \gamma_\mu P_L q)$$



$\mu \rightarrow e$ conversion



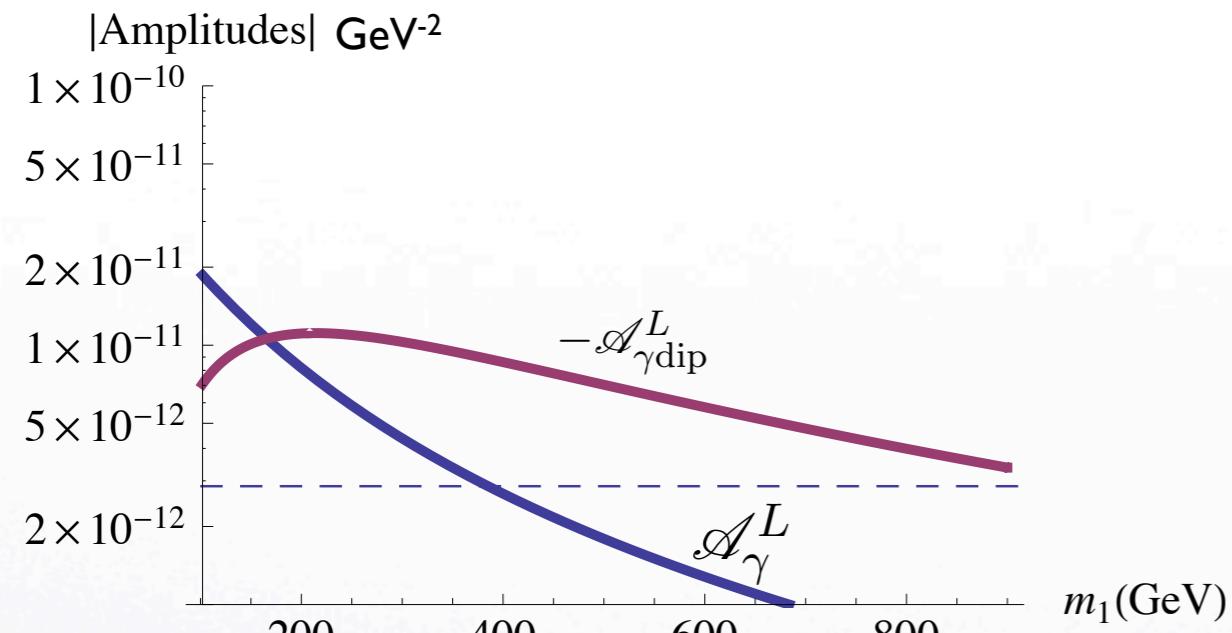
- No chirality flip on quark line, $\bar{q}\gamma^\mu q$, $\bar{q}\gamma^\mu\gamma^5 q$
 - Considering only spin independent term, $\bar{q}\gamma^\mu q$ remains
- \Rightarrow No sea quark contributions, only u and d quarks



Maximal mixing

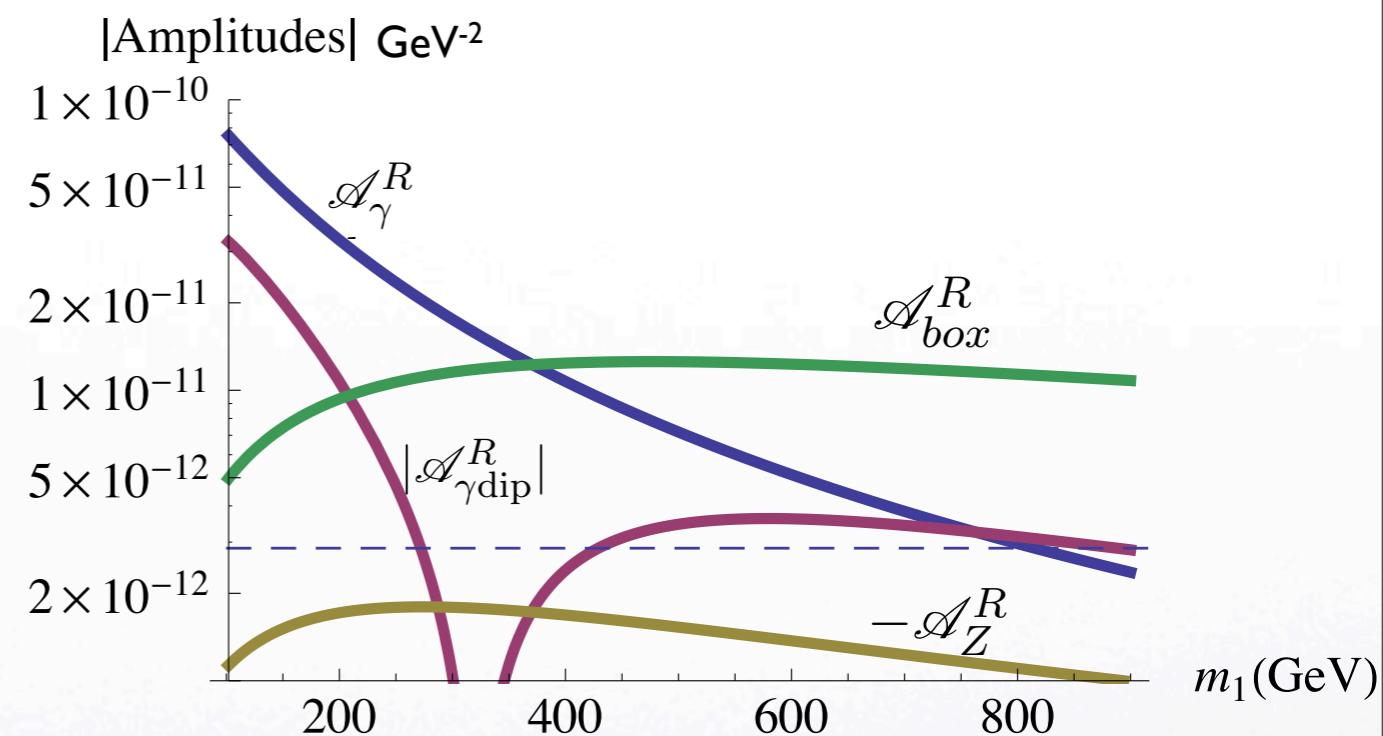


(d) $M_1 = \mu_d = 200$ GeV



LH

(d) $M_1 = \mu_d = 200$ GeV

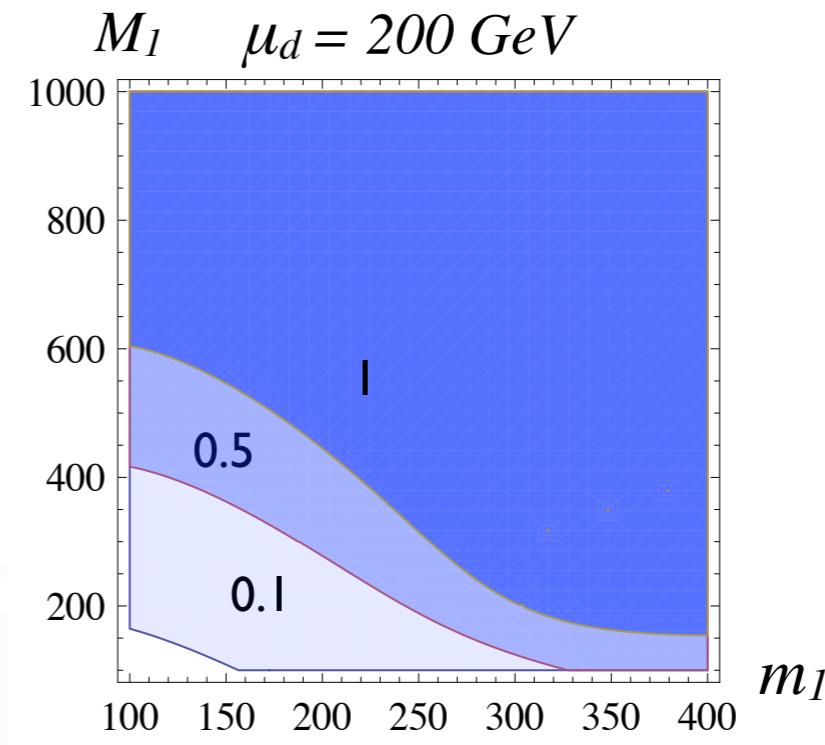
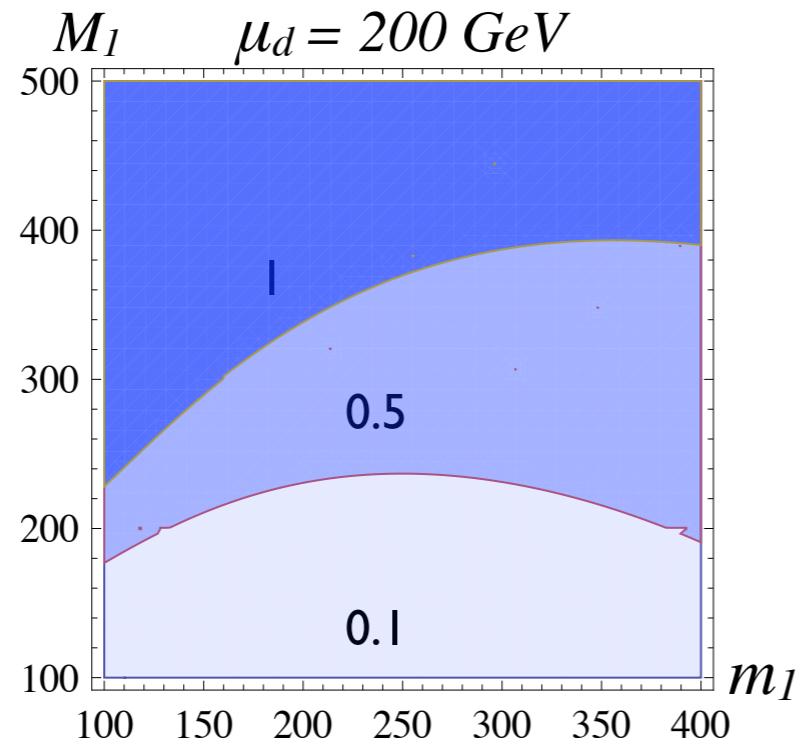


RH

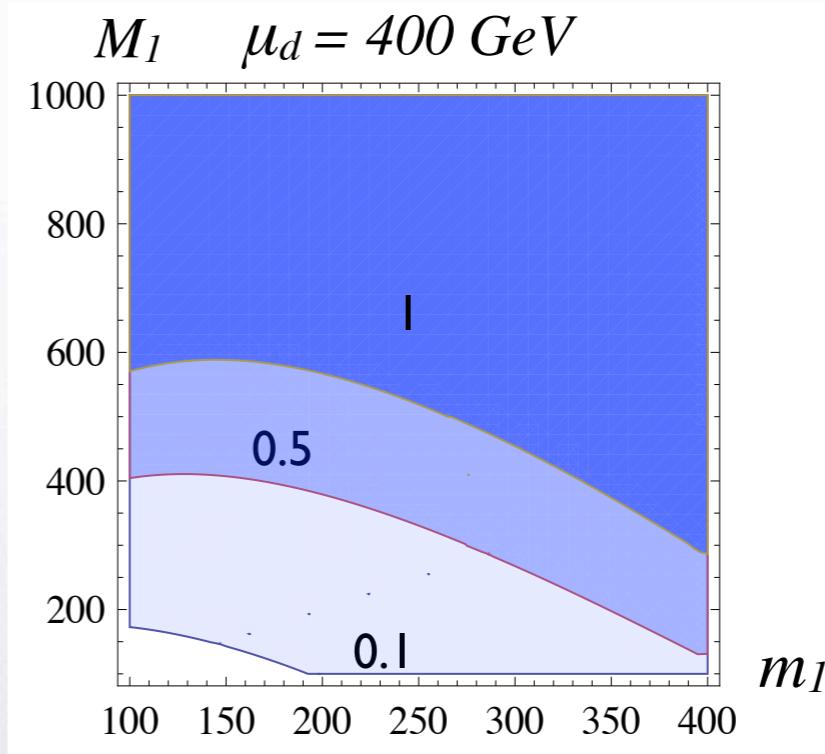
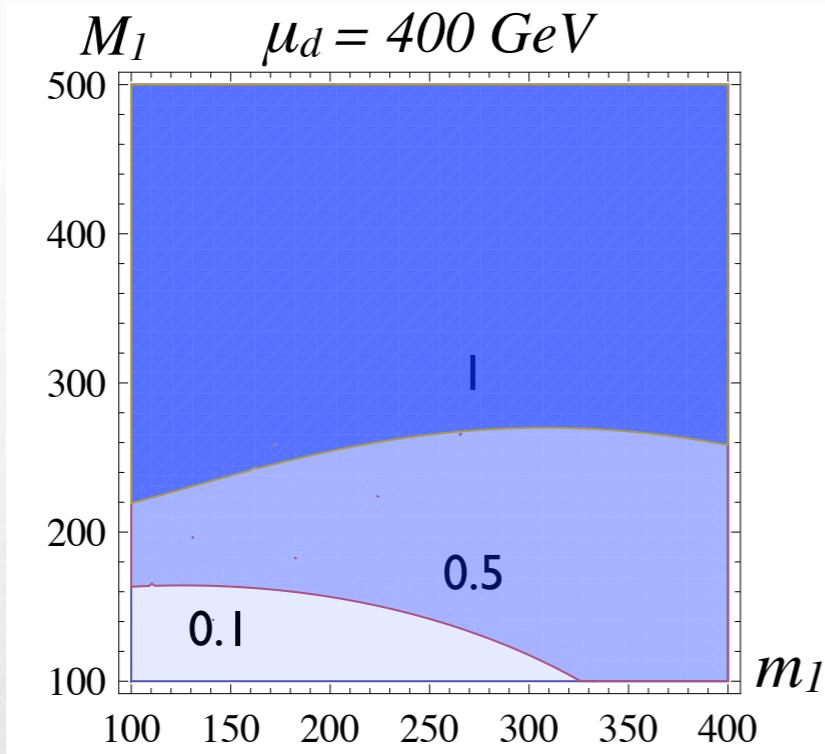
$$\Gamma_{\mu \rightarrow e} = 4m_\mu^5 e^4 |\mathcal{A}_{\gamma\text{dip}}^L + \mathcal{A}_\gamma^R + \mathcal{A}_{\text{box}}^R + \mathcal{A}_Z^R|^2 + (L \leftrightarrow R)$$



LH



LH

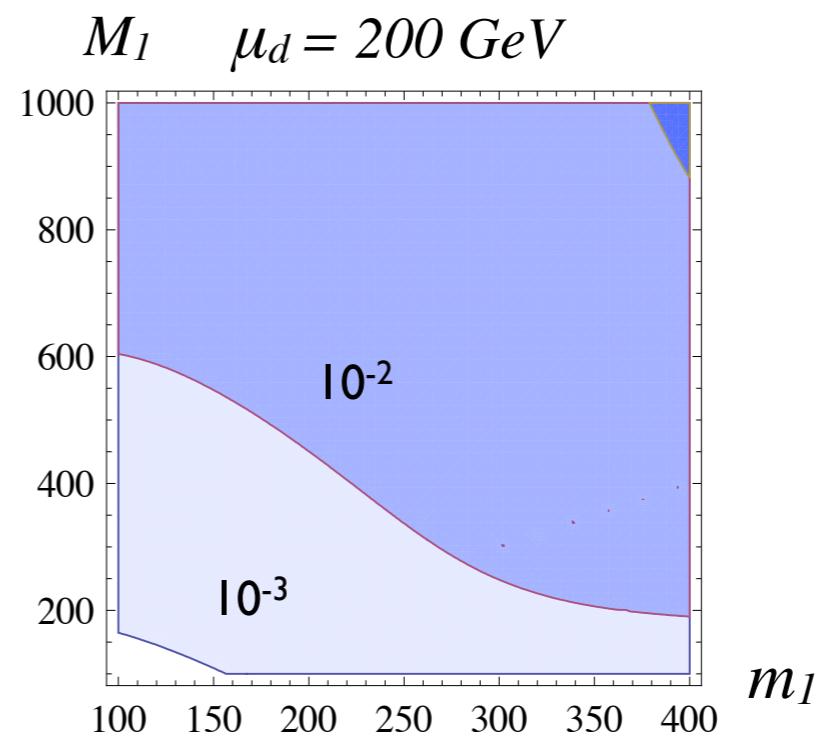
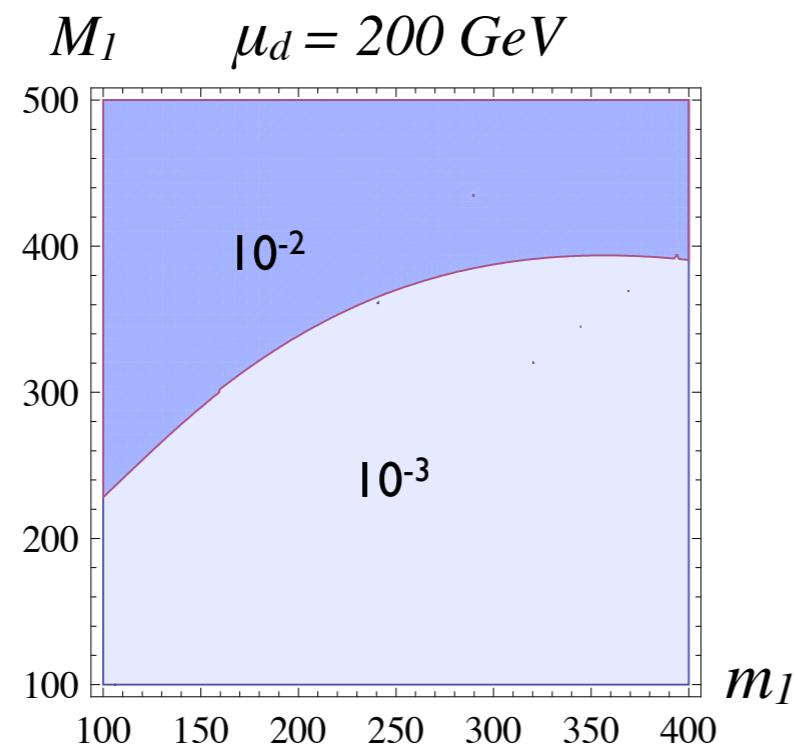


RH

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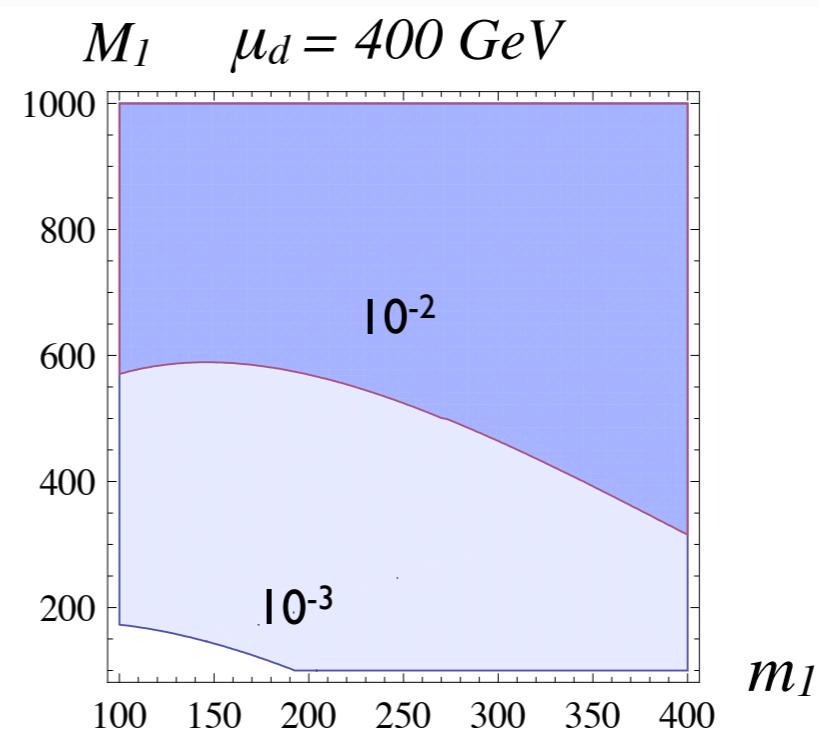
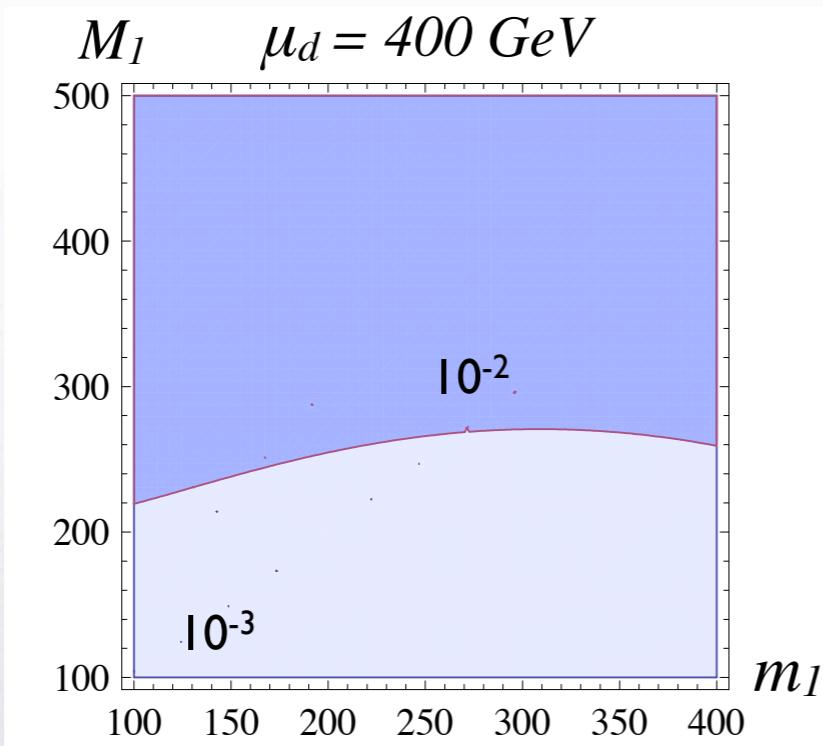


LH



RH

LH



RH

Fok